

---

## Lecture 2: Invariants, cross-section, Feynman diagrams

- Last lecture introduced:
    - 4-vectors
    - 4-vector multiplication
    - Lorentz transformations
    - Introduced : transverse mass, transverse momentum , rapidity, pseudo-rapidity
    - Rapidity transformation with Lorentz boost
    - Showed that  $m_T$ ,  $m$  ,  $p_T$  are Lorentz invariant
- 

---

## Today's lecture

- Now that we know the appropriate language to describe the kinematics, we want to move to a more exciting things like collisions (i.e. interactions) between particles
  - We start with AA collisions in the initial state and detect particles:  $\pi, K, p, \gamma, e^+, e^-, \mu^+, \mu^-$  in the final state
  - We can reconstruct particles that have decayed:  $\Lambda, J/\psi, \phi, \Omega \dots$
  - We want to know if AA is just a superposition of NN or if we have created a QGP medium that affects the particle production
  - => we need to understand NN collisions first
-

---

## Two main processes to consider

- Scattering: 2->2 and 2-> many
  - Decay: 1->2 ; 1-> several
  - We need to describe the production rate of particles in the final state in a frame-independent way
  - Cross-section : has dimension of area .  
Measured in barns =  $10^{-24} \text{ cm}^2$
- 

---

## Cross-section

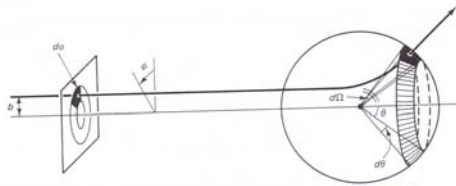


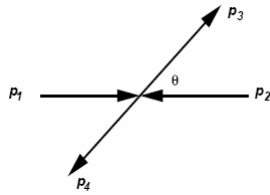
Figure 6.3 Particle incident in area  $d\sigma$  scatters into solid angle  $d\Omega$ .

$$L = \frac{fnN_1N_2}{A}$$

$$\frac{d\sigma}{d\Omega} = \frac{1}{L} \frac{d^2N}{d\Omega dt}$$

- Cross-section: particles coming through area  $d\sigma$  scatter in  $d\Omega$
  - Luminosity:
    - $f$  is the revolution frequency
    - $n$  is the number of bunches in one beam in the storage ring.
    - $N_i$  is the number of particles in each bunch
    - $A$  is the cross section of the beam.
    - RHIC design luminosity:  $2 \times 10^{26} \text{ cm}^{-2} \text{ s}^{-1}$
    - Currently:  $\sim 5 \times 10^{26} \text{ cm}^{-2} \text{ s}^{-1}$
-

## 2-> 2 scattering and invariants



- Particle (1) and (2) in the initial state
- Particle (3) and (4) in the final state
- $p_1, p_2, p_3, p_4$  – 4-momenta

Mandelstam invariants:

$$\begin{aligned} S &= (\mathbf{p}_1 + \mathbf{p}_2)^2 = (\mathbf{p}_3 + \mathbf{p}_4)^2 \\ T &= (\mathbf{p}_1 - \mathbf{p}_3)^2 = (\mathbf{p}_2 - \mathbf{p}_4)^2 \\ U &= (\mathbf{p}_1 - \mathbf{p}_4)^2 = (\mathbf{p}_2 - \mathbf{p}_3)^2 \end{aligned}$$

## Do this on the board or for homework

- Show that  $S$  is Lorentz invariant
- Show that  $\text{sqrt}(S) = E_{\text{cm}}$
- Find the relationship between the beam momentum  $p_1$  in a fixed target experiment and the center of mass momentum
- Show that :  
 $S+T+U = \sum_i m_i^2$
- Show that  $d^3p/E$  is invariant

## The Golden rule

- Cross section and decay rates are expressed in terms of:
  - The amplitude for the process:  $M$  ( or matrix element between the initial and final state).
    - Contains all dynamical information
    - Calculated by evaluating appropriate Feynman diagrams using the Feynman rules for the appropriate interaction in question
  - The available phase space ( or density of final states) – contains the kinematical information
  - Example: decay of a heavy particle into many light secondaries – large phase space – many different ways to apportion the available energy. Neutron decay – small phase space, because  $m_p$  is close to  $m_n$

$$\text{transition rate} = \frac{2\pi}{\hbar} |M|^2 \times \text{phase space}$$

## Golden rule (II)

- 1->2+3+....+n

$$dLips = (2\pi)^4 \delta^4(p_1 - p_2 - p_3 - \dots - p_n) \frac{d^3 p_2}{(2\pi)^3 2E_2} \frac{d^3 p_3}{(2\pi)^3 2E_3} \dots \frac{d^3 p_n}{(2\pi)^3 2E_n}$$

$$d\Gamma = |M|^2 \frac{dLips}{2m_1}$$

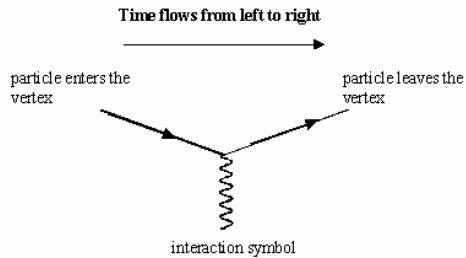
- 1+2->3+4+..n

$$dLips = (2\pi)^4 \delta^4(p_1 + p_2 - p_3 - p_4 \dots - p_n) \frac{d^3 p_3}{(2\pi)^3 2E_3} \frac{d^3 p_4}{(2\pi)^3 2E_4} \dots \frac{d^3 p_n}{(2\pi)^3 2E_n}$$

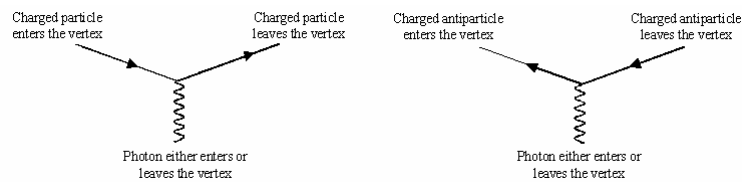
$$d\sigma = |M|^2 \frac{dLips}{\text{initial flux}}$$

## What goes into M ?

- Here we need to consider the interaction
  - Note: a vertex is simply a symbol, it does not represent tracks of particles in space and it is not a space-time diagram
  - At each vertex, energy and momentum is conserved
  - At each vertex, we put the coupling constant , i.e. -> the strength of the interaction

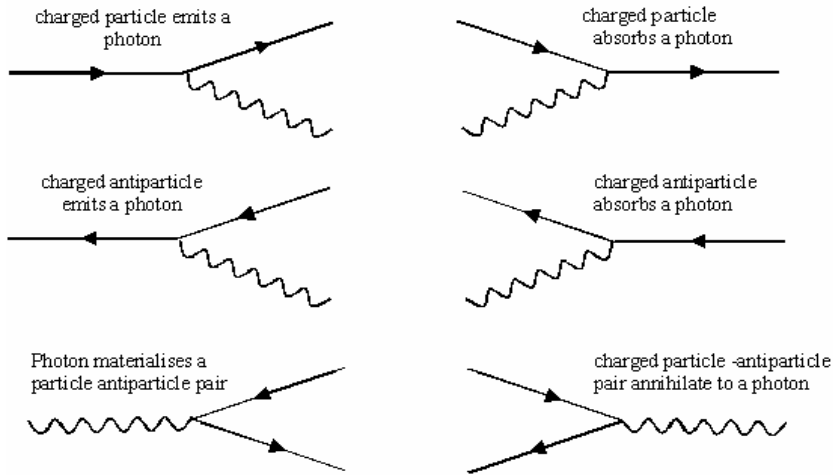


## Electromagnetic interaction



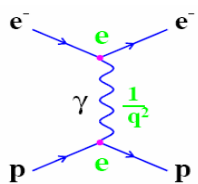
- The coupling is  $\sqrt{\alpha}$  ,  $\alpha = e^2/4\pi \sim 1/137$

# QED processes: rotate the legs around

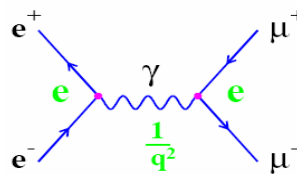


# QED - continued

electron-proton scattering

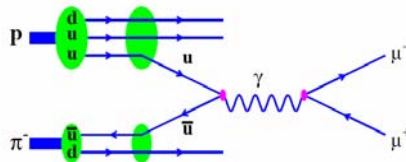


electron-positron annihilation



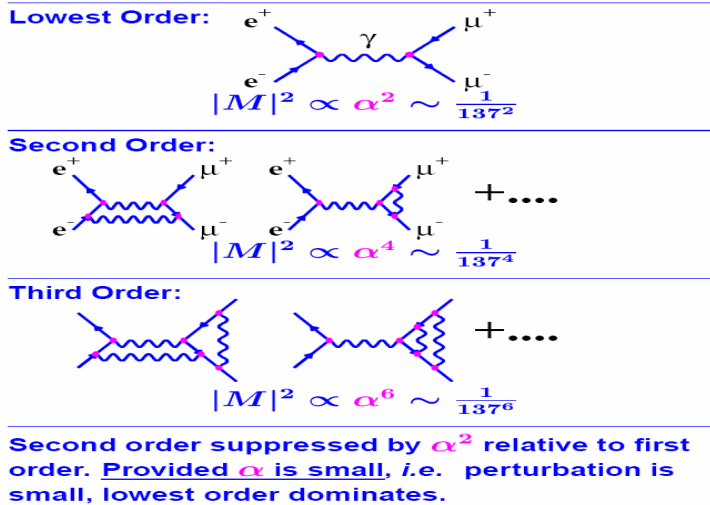
$$\pi^- p \rightarrow \mu^+ \mu^- + \text{hadrons}$$

Quark-anti-quark annihilation: Drell-Yan

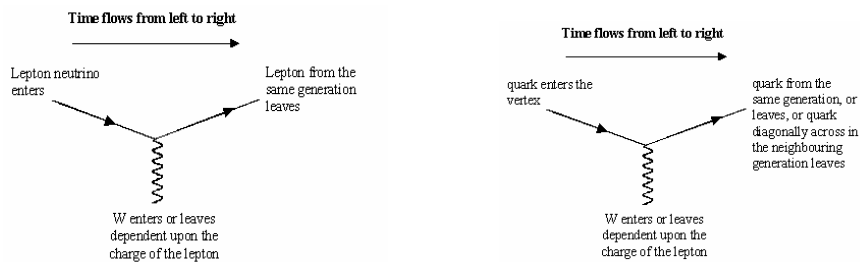


—  $M \sim$  is expressed in the same way  $\sim e^2/q^2$  However, the four-momentum transfer is very different —

# Perturbative expansion

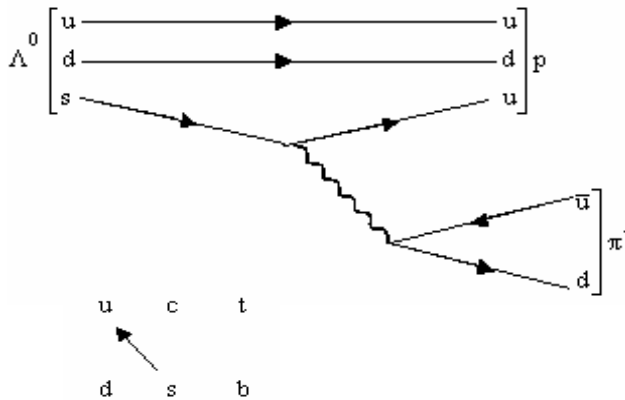


# Weak vertices

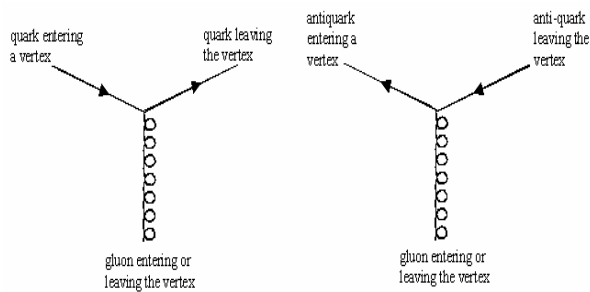


- We could also exchange a Z boson: any quark or lepton doesn't change identity
- Additional vertices: W couples to Z, W and to photon
- NOTE: the mass of the gauge boson goes in the denominator of the propagator. That's what makes the interaction "weak". The coupling constant is actually larger Than the EM coupling constant.

## Weak decay example: quarks change generation



## Strong interaction

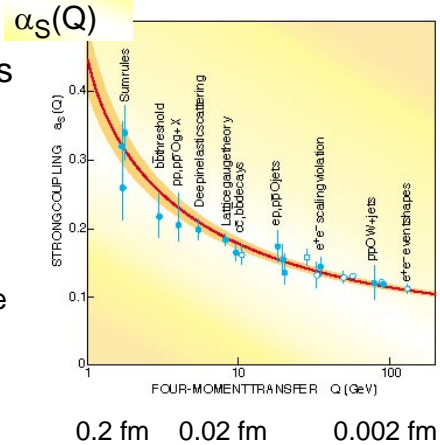


Gluons carry color+anti-color. They couple to themselves and to quarks.  
 The interaction is strong due to the coupling constant being  $\sim 1$ .  
 Higher order diagrams contribute more  $\otimes$ . Can't do a perturbative expansion!  
 We get one lucky break. The coupling constant is NOT a constant !



## QCD's Key Feature

- Hadron masses  $\sim 1$  GeV
- Hadron sizes  $\sim 10^{-15}$  meters (aka 1 femtometer or 1 fermi = 1 fm )
- Planck's constant  $\hbar c = 0.2$  GeV-fm
  - ➔  $1 \text{ fm}^{-1} \Leftrightarrow 200 \text{ MeV}$
  - ➔  $200 \text{ MeV} \sim$  characteristic scale associated with confinement
- These are reflected in the “running coupling constant” of QCD
- Can do pQCD when  $Q^2$  is large “hard scattering”



## Cross-sections in pp collisions

- Total = elastic + inelastic
- Elastic  $\sim 10$  mb
- Inelastic  $\sim 40$  mb at RHIC energies
- Energy dependence : parameterized from data

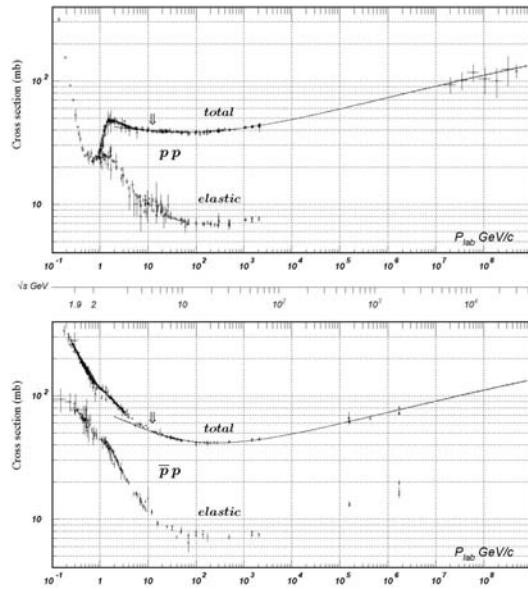


Figure 40.11: Total and elastic cross sections for  $pp$  and  $\bar{p}p$  collisions as a function of laboratory beam momentum and total center-of-mass energy. Corresponding computer-readable data files may be found at <http://pdg.lbl.gov/xsect/contents.html> (Courtesy of the COMPAS group, IHEP, Protvino, August 2005)