

# Lecture 08: transverse energy, nuclear stopping and energy density

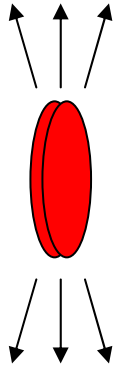
- Last lecture:
  - Nuclear geometry, centrality,  $N_{\text{part}}$ ,  $N_{\text{coll}}$
  - Multiplicity:
    - Energy dependence
    - Pseudo-rapidity dependence and scaling
      - at mid-rapidity
      - Total multiplicity
      - fragmentation region
- Today: nuclear stopping, transverse energy and energy density – i.e. get a more detailed picture of particle production and energy flow

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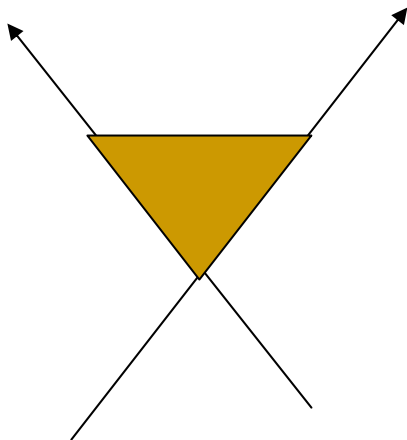
## Consider 2 pictures of the AA collisions

- Colliding nuclei stop and convert all of their kinetic energy into particles (Landau)
  - Colliding nuclei are transparent to each other. They “pass-through” stretching strings between the colliding partons. As the strings break – particles are produced. Only a fraction of the initial kinetic energy is lost and available for particle production. (Bjorken)
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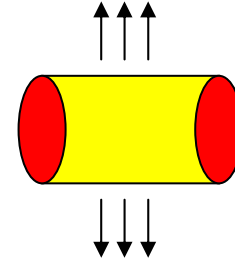
# Landau vs Bjorken picture



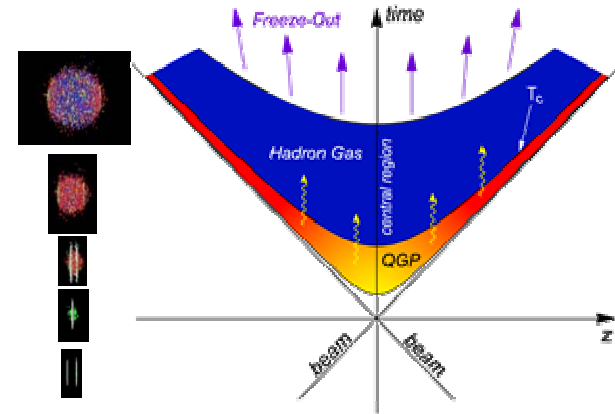
Landau



Gaussian rapidity distribution.



Bjorken



Rapidity plateau, or "boost invariance".

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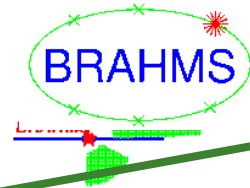
Let's measure where the initial baryons go – then we'll know which picture is correct

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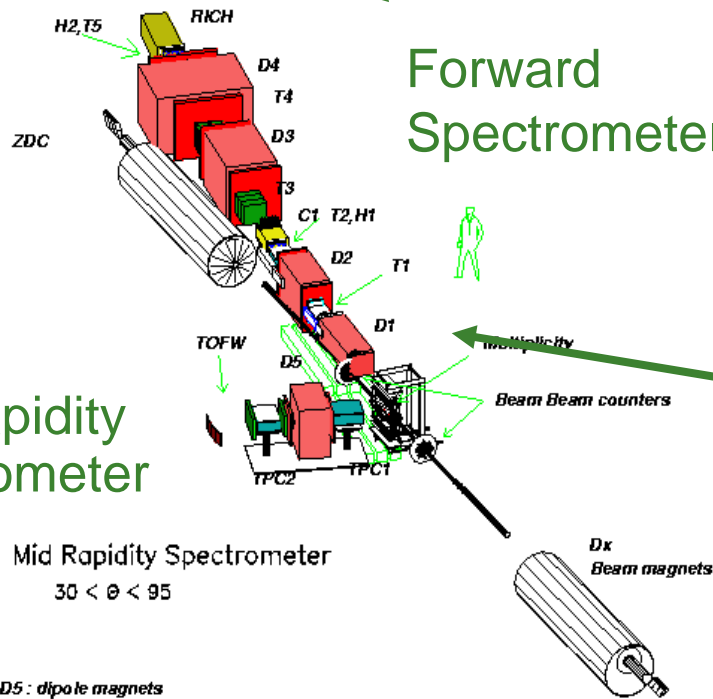
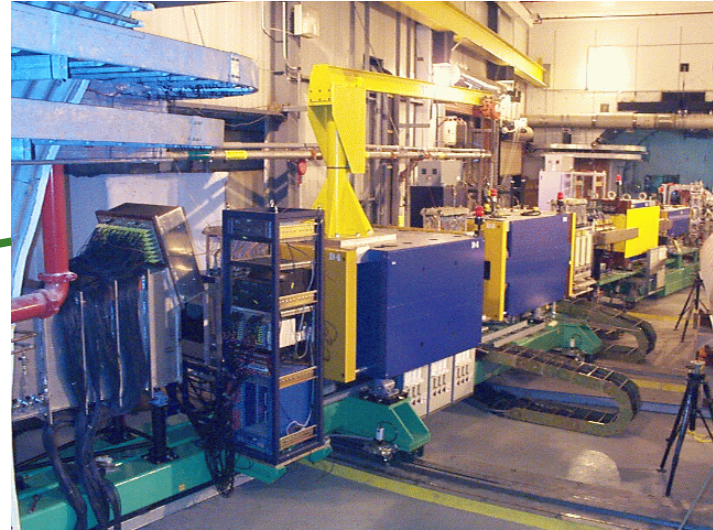
# The Brahms detector

Two Arm Spectrometer  
with Variable Setting

Good Particle ID

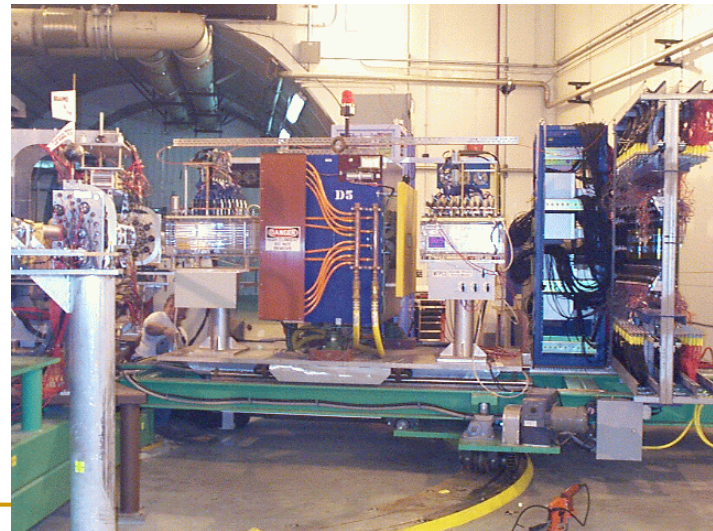


Forward  
Spectrometer



Mid-Rapidity  
Spectrometer

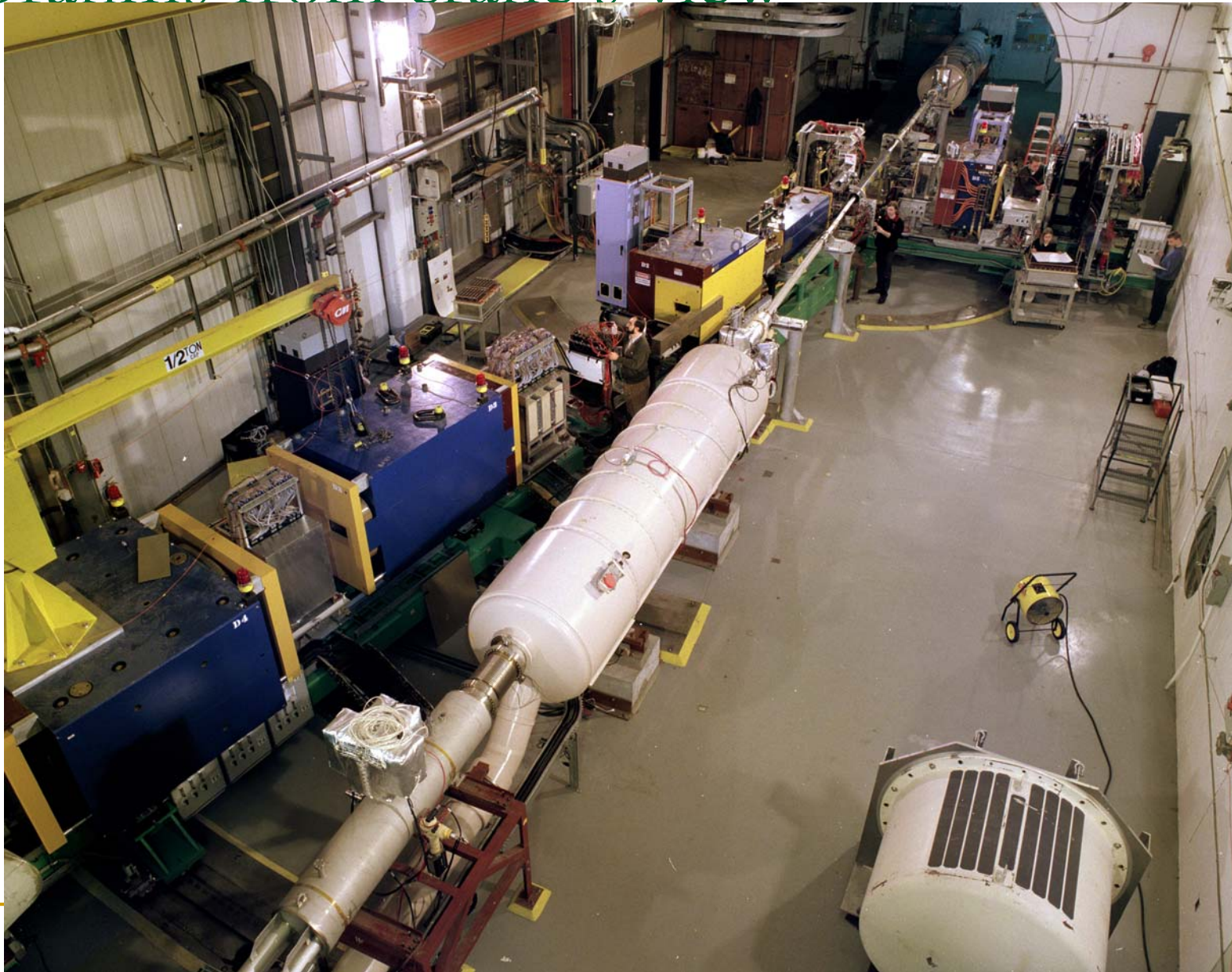
Mid Rapidity Spectrometer  
 $30 < \theta < 95$



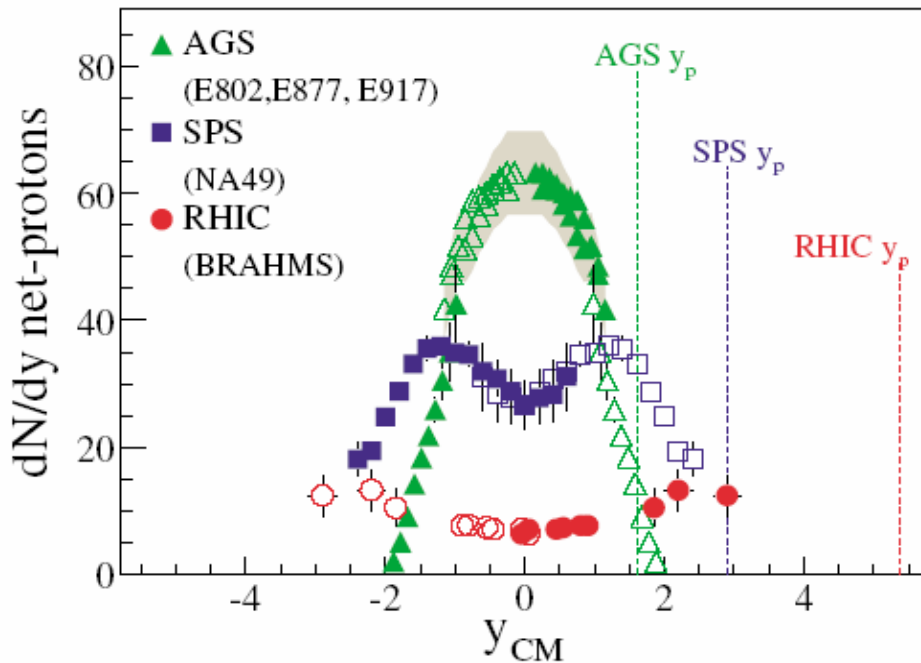
D1, D2, D3, D4, D5 : dipole magnets  
T1, T2, T3, T4, T5, TPC1 TPC2: tracking detectors  
H1, H2, TOFW : Time-of-flight detectors  
RICH, GASC : Cherenkov detectors



# Brahms from crane's view



# Net baryon (these are the transported particles) rapidity distributions



BRAHMS Collaboration (I. G. Bearden et al.)

["Nuclear Stopping in Au+Au Collisions at  \$\sqrt{s\_{NN}}=200\$  GeV"](#)

**Phys. Rev. Lett. 93, 102301 (2004)**

FIG. 3 (color online). The net-proton rapidity distribution at AGS [8,21,22] (Au + Au at  $\sqrt{s_{NN}} = 5$  GeV), SPS [23] (Pb + Pb at  $\sqrt{s_{NN}} = 17$  GeV), and this measurement ( $\sqrt{s_{NN}} = 200$  GeV). The data are all from the top 5% most central collisions and the errors are both statistical and systematic (the light gray band shows the 10% overall normalization uncertainty on the E802 points, but not the 15% for E917). The data have been symmetrized. For RHIC data black points are measured and gray points are symmetrized, while the opposite is true for AGS and SPS data (for clarity). At AGS weak decay corrections are negligible and at SPS they have been applied.

# Net proton rapidity distributions

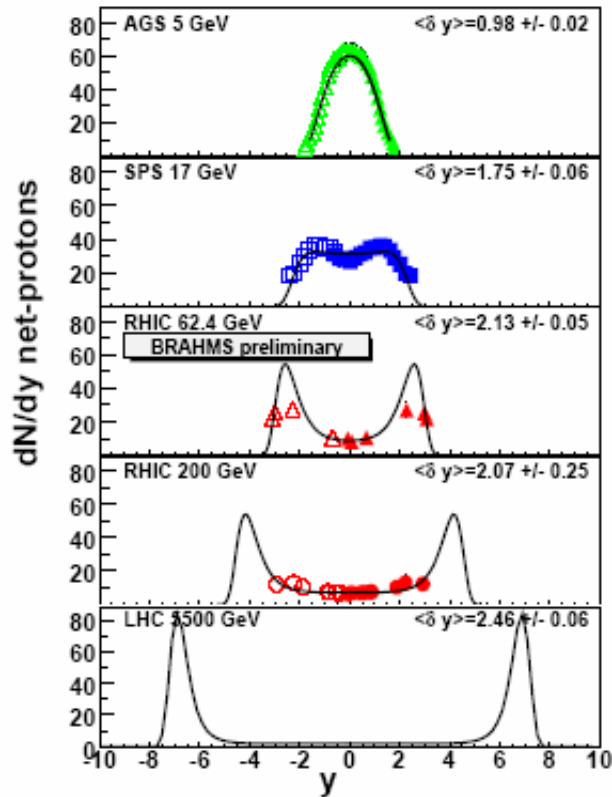


Fig. 3.  $\frac{dN}{dy}$  of protons from AGS<sup>6</sup>, SPS<sup>9</sup> and RHIC<sup>5</sup>. Errors include systematic errors. Bottom panel shows extrapolation to LHC.

Since baryon number is conserved,  
We can extrapolate the net baryon  
distributions in the unmeasured  
region.

Then measure the average rapidity  
loss to quantify the stopping:

$$\langle \delta y \rangle = y_p - \langle y \rangle$$

Here,  $y_p$  is the rapidity of the  
incoming projectile and  $\langle y \rangle$  is the  
mean net-baryon rapidity after  
the collision:

$$\langle y \rangle = \frac{2}{N_{\text{part}}} \int_0^{y_p} y \frac{dN_{(B-\bar{B})}(y)}{dy} dy.$$



# Nuclear stopping from AGS to LHC

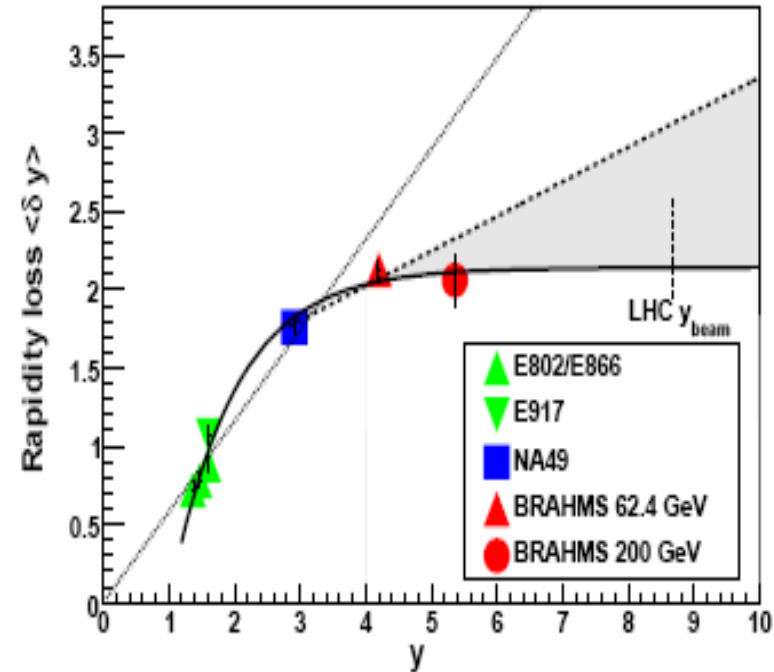
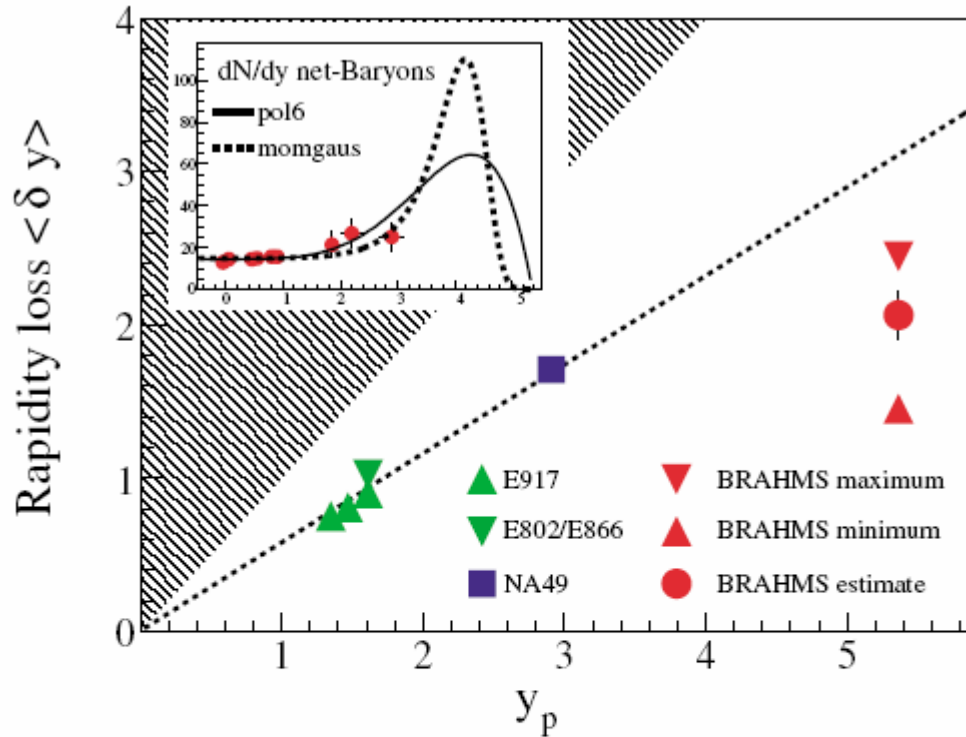
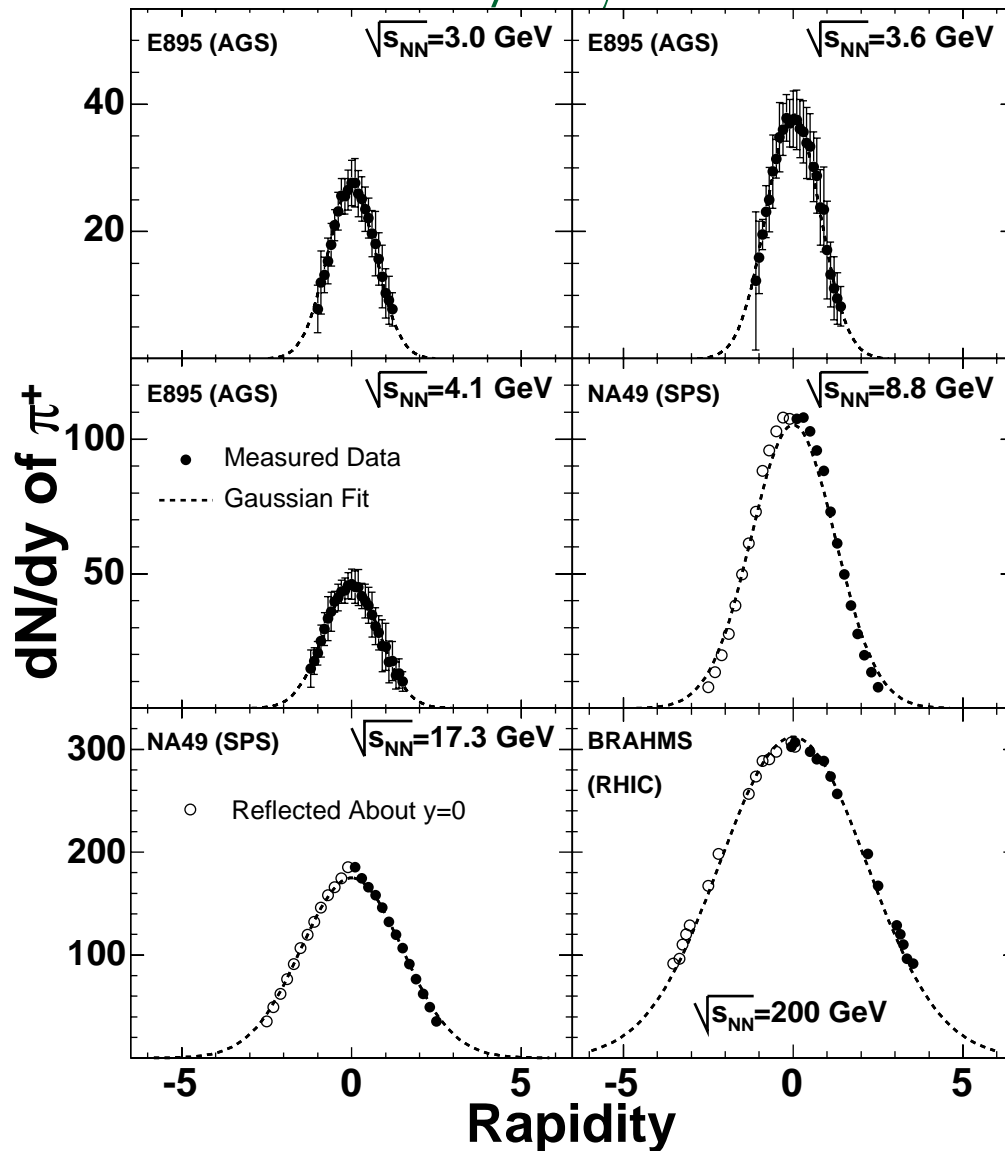


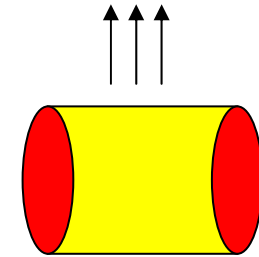
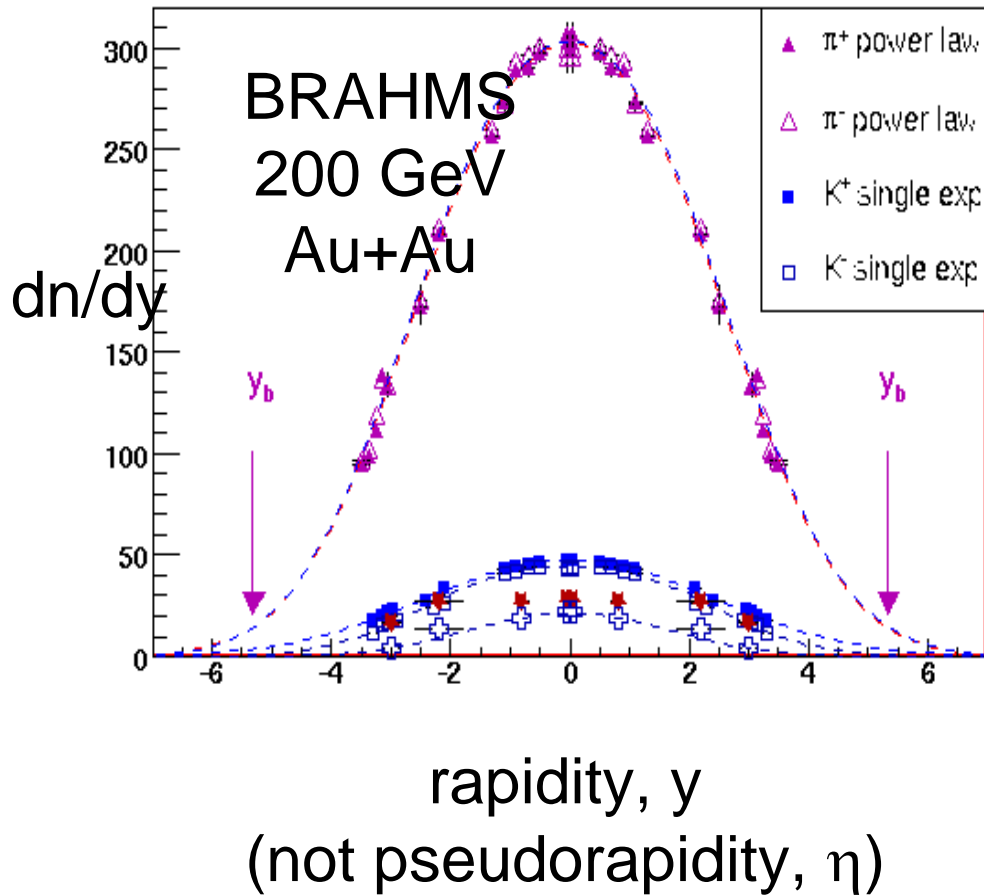
FIG. 4 (color online). The inset plot shows the extrapolated net-baryon distribution (data points) with fits (represented by the curves) to the data; see text for details. The full figure shows the rapidity loss, obtained using Eq. (1), as a function of projectile rapidity (in the CM). The hatched area indicates the unphysical region, and the dashed line shows the phenomenological scaling  $\langle \delta y \rangle = 0.58 y_p$ . The data from lower energy are from [6,8].

# Look at $dN/dy$ – no boost invariance



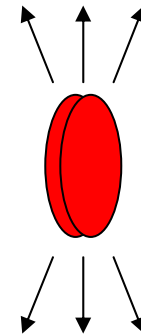
- No plateau is present in the  $dN/dy$  distributions of pions

# Boost invariant?



Bjorken

The “truth” is in-between.



Landau

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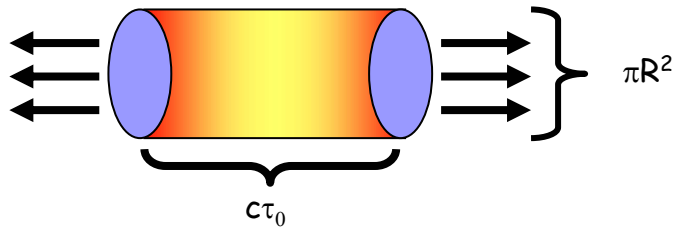
Now go on and estimate the energy density in the collisions

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This will let us find out if we have exceeded the critical energy density for the phase transition.

energy density =  $\Sigma E / \text{volume}$

Colliding system expands:



Energy  $\perp$  to  
beam direction  $\downarrow$

$$\mathcal{E}_{Bj} = \frac{1}{\pi R^2} \frac{1}{c \tau_0} \left( \frac{dE_T}{dy} \right)$$

per unit  $\nearrow$   
velocity  $\parallel$  to beam

$$\mathcal{E}_{Bj} = \frac{1}{\pi R^2} \frac{1}{c \tau_0} \frac{dN}{dy} \langle m_T \rangle$$

**Question:**

What is the relevant time during the collision at which we need to calculate the energy density ?

**Note:** Lattice calculations predict critical energy density  $\varepsilon_c \sim 0.3\text{-}1.0 \text{ GeV}/\text{fm}^3$ .

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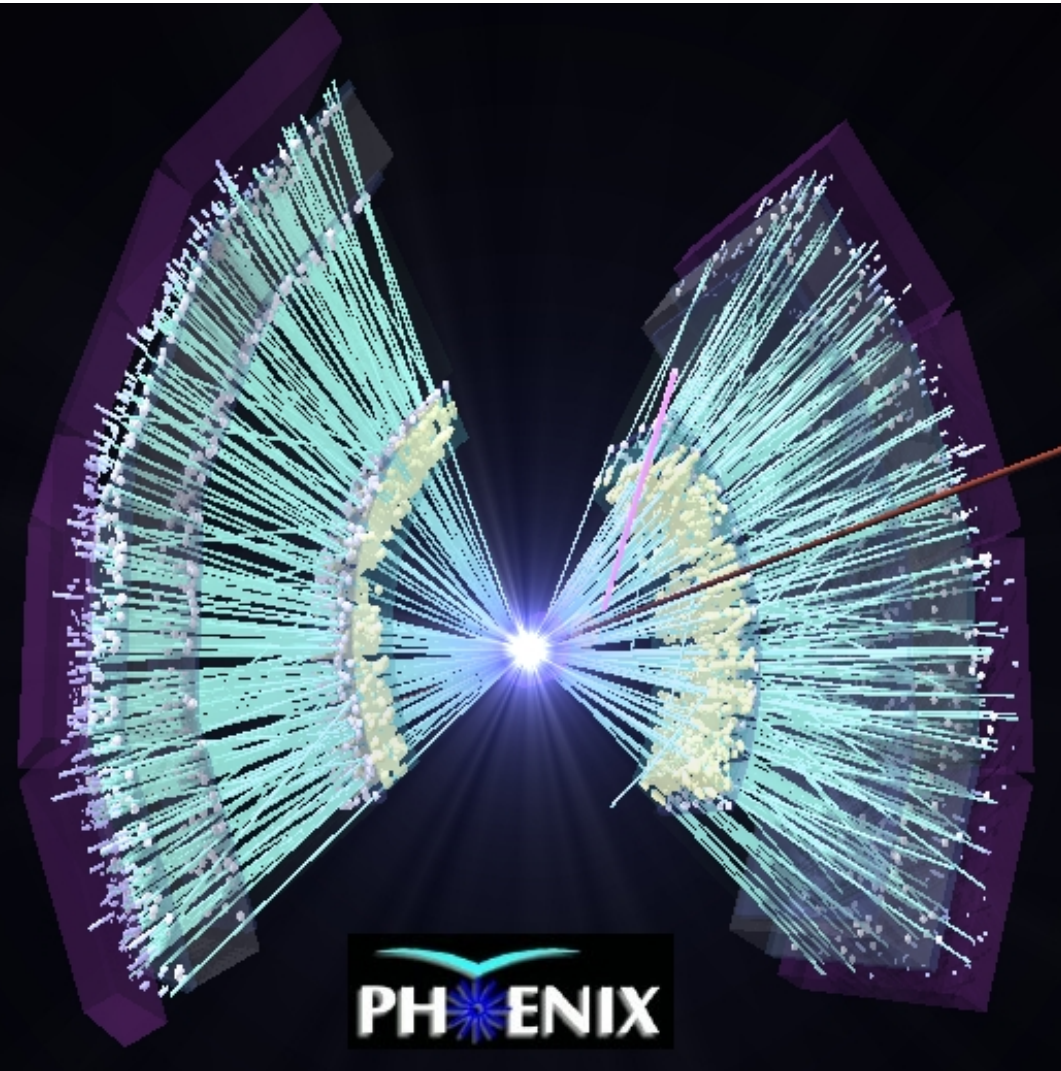
# Estimating energy density

- Time less than the time needed for the nuclei to pass through doesn't make sense, because  $\varepsilon$  becomes unphysically large trivially (just by overlap mass), so we need:

$$\tau \gg \frac{2R}{\gamma}$$

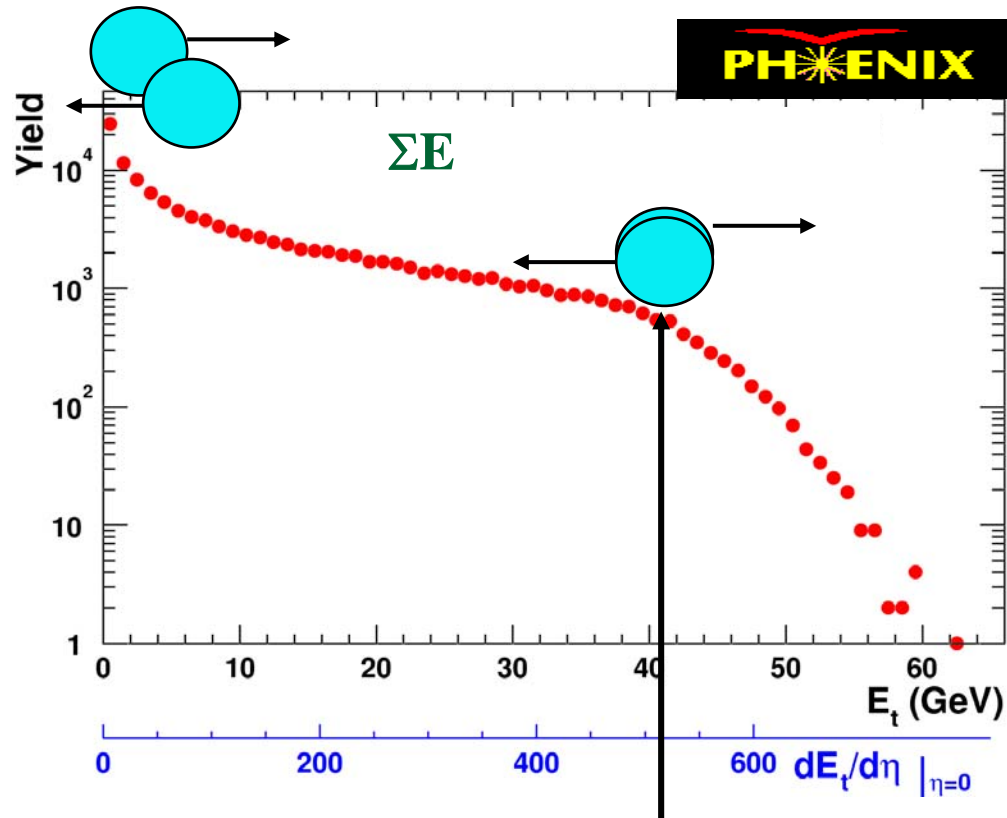
- For RHIC full energy  $\gamma = 106 \Rightarrow \tau_{\text{pass through}} \sim 0.13 \text{ fm}/c$
- We need to consider “formed” or secondary particles – following Bjorken (PRD 27 (1983) 140) –  $\tau_{\text{form}} \sim 1 \text{ fm}/c$
- So, let's measure the transverse energy and get the energy density
- We also need to know how to define centrality (to get the volume), and we discussed this already

# How do we measure $E_T$ ?

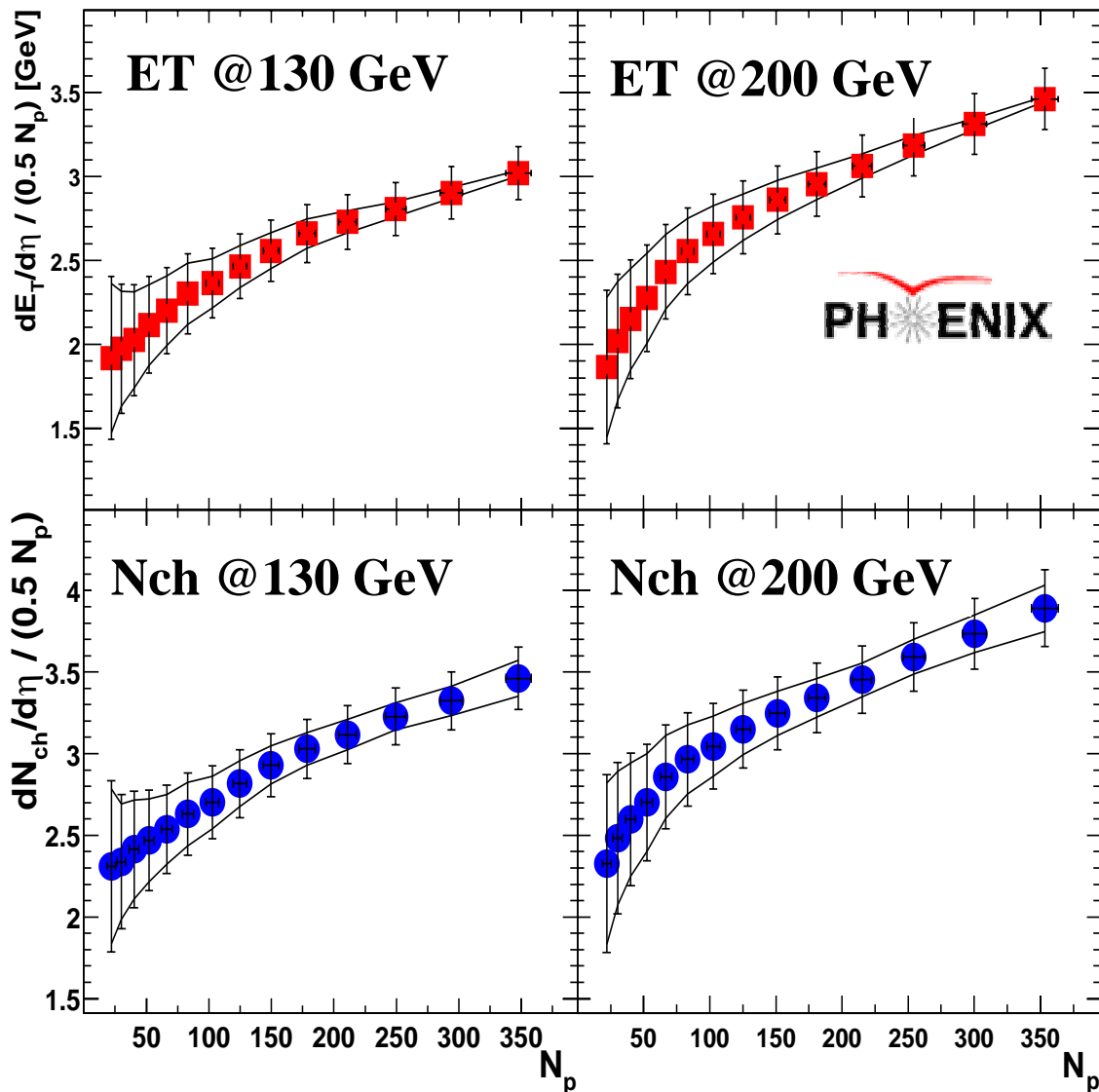


**We measure the particles coming out, so add up their energy.  
Put a calorimeter detector (measures energy) – sum it up for all particles.**

# Transverse energy at mid-rapidity



# Results for $E_T$ and $N_{ch}$ : centrality dependence



**Stat. errors**

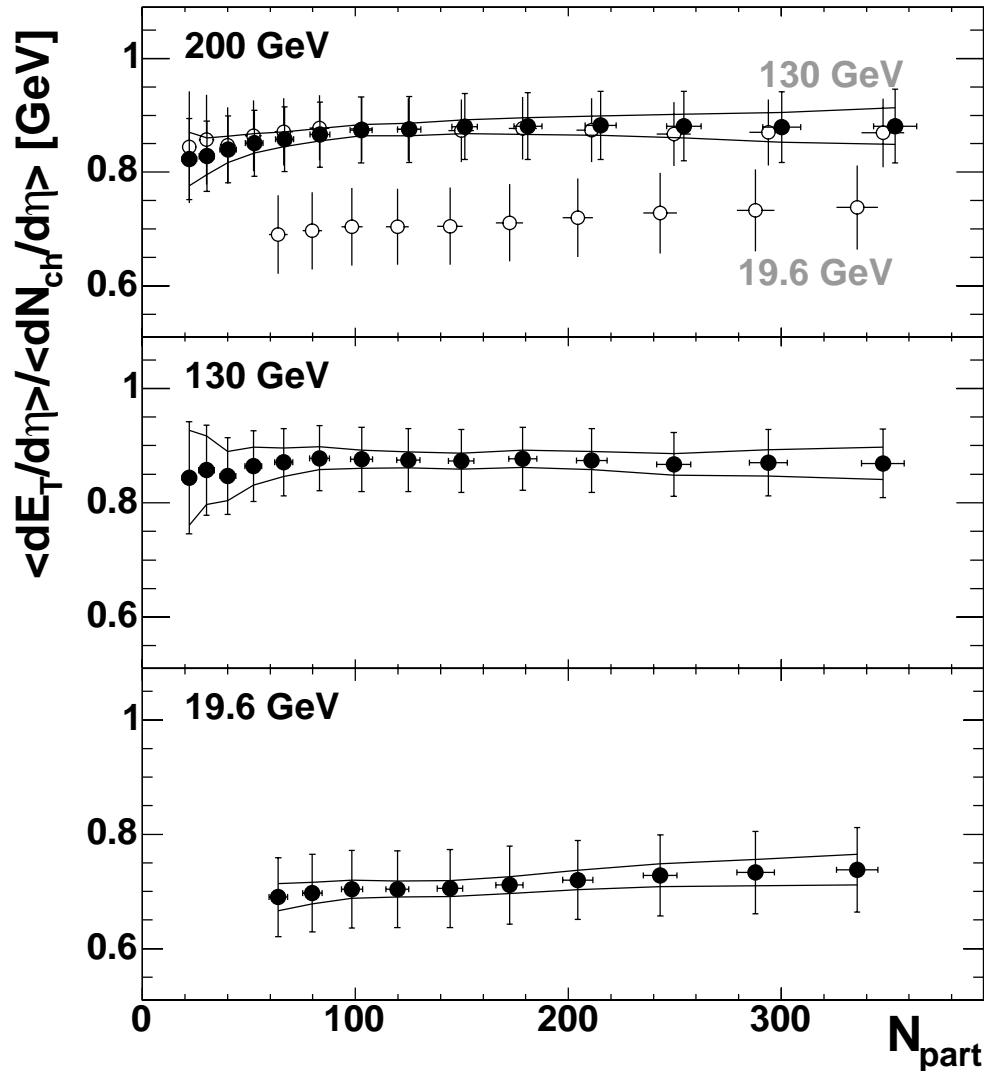
Negligible

**Syst. errors**

Band: possible  
common tilt

Bars: total syst. error

# $E_T / N_{ch}$



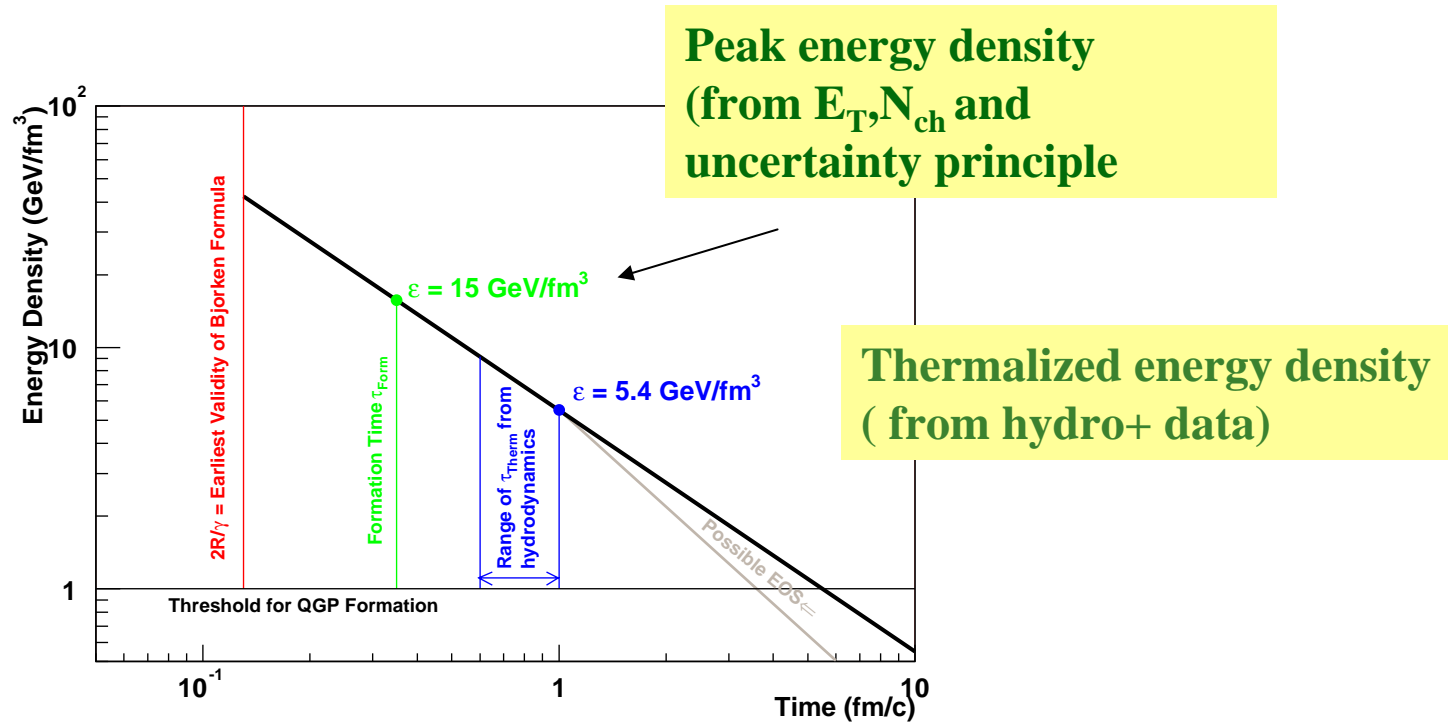
- Remarkably,  $dE_T/d\eta / dN_{ch}/d\eta$  does not change much with  $\sqrt{s}$
- The extra energy goes into particle production
- From  $dE_T/d\eta / dN_{ch}/d\eta = 0.85$  GeV after converting to  $dN/dy$ , we get  $\langle m_T \rangle \sim 0.57$  GeV
- If we assume that

$$\tau_{form} \approx \hbar / \langle m_T \rangle$$

We get  $\tau_{form} \sim 0.35$  fm/c . This is smaller than the “nominal”, but larger than  $2R/\gamma$



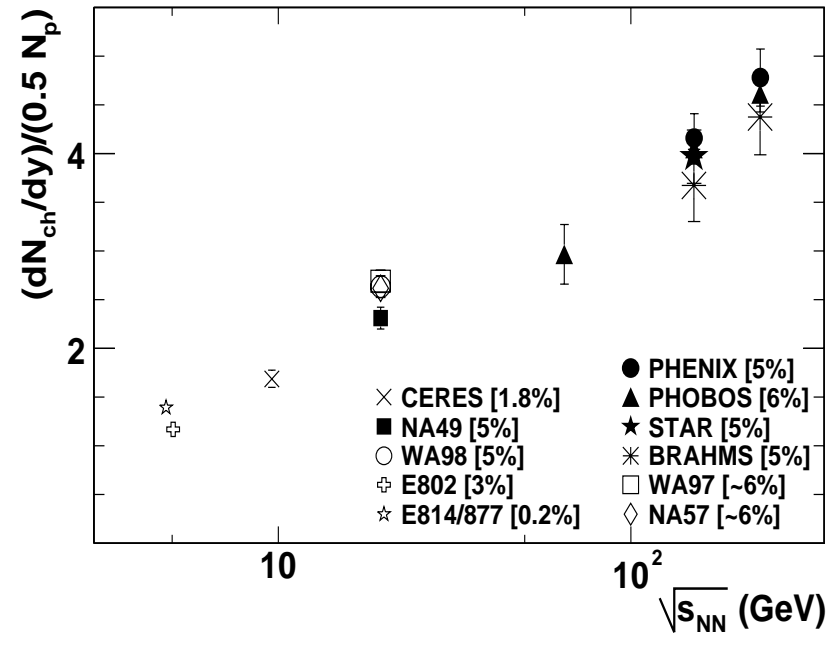
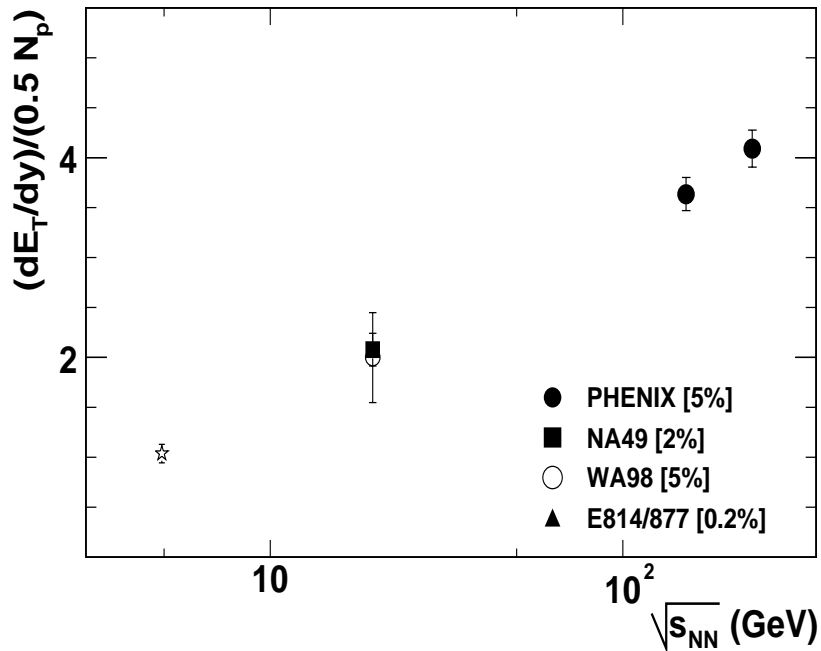
# Now ... estimate $\epsilon$



- With the “nominal”  $\tau_{\text{form}} = 1 \text{ fm}/c$ ,
  - $\epsilon \geq 5.5 \text{ GeV}/\text{fm}^3$  (200 GeV Au+Au)
- With the uncertainty principle limit:  $\tau_{\text{form}} = 0.35 \text{ fm}/c$ ,
  - $\epsilon \geq 15 \text{ GeV}/\text{fm}^3$  (200 GeV Au+Au)

→ *well above predicted transition!*

# Energy dependence of energy density



- Both  $dN_{ch}/d\eta$  and  $dE_T/d\eta$  show logarithmic growth with  $\sqrt{s_{NN}}$
- At LHC expect ~ factor 20 increase in  $\epsilon$ .  
 $dN_{ch}/d\eta \sim 1200$

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Extra slide on  $E_T$

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# Transverse Energy Measurements

**Convention:**

$$E_T = \sum E_i \sin \theta_i$$

$E_i = E_i^{tot} - m_N$  for baryons  
 $E_i = E_i^{tot} + m_N$  for antibaryons  
 $E_i = E_i^{tot}$  for others

EMCal is “almost” hadronic calorimeter  
(depth 18 radiation lengths, or 0.85  
interaction lengths)

$$E_{EMC} = 1.0 \cdot E_{tot} \text{ for } \gamma, \pi^0$$
$$E_{EMC} = 0.7 \cdot E_{tot} \text{ for } \pi^\pm$$

**$E_{EMC} \rightarrow E_T$  transformation:**

$$E_T = 1.23 \cdot E_{EMC}$$

Do a MC simulation with realistic  
particle composition

## EMCal absolute energy calibration

MIP peak

$E/p$  matching peak for  $e^\pm$

$\pi^0$  mass peak

